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**SPATIAL DISPERSION PROPERTIES FROM DFPT:
DYNAMICAL QUADRUPOLES AND
FLEXOELECTRIC TENSOR**

The background of the slide is a dark blue gradient. Overlaid on this are several glowing, ethereal blue lines that form abstract, flowing shapes. These lines vary in thickness and brightness, creating a sense of movement and depth. Some lines are more prominent, while others are fainter, blending into the background. The overall effect is reminiscent of a complex, dynamic system or perhaps a visualization of physical phenomena related to the text.

*9th ABINIT Developers Workshop
Louvain-la-Neuve, May 2019*

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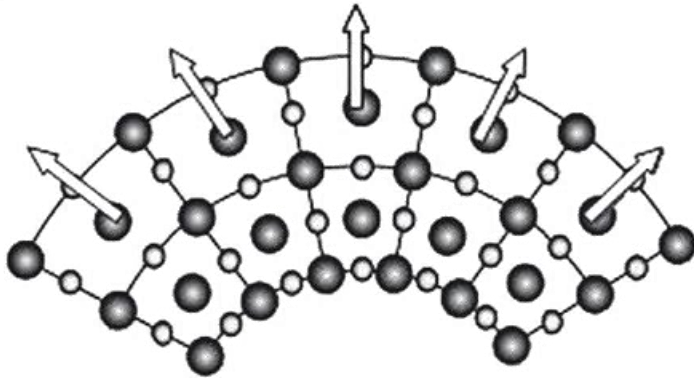
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Flexoelectricity



$$P_\alpha = \mu_{\alpha\beta,\gamma\delta} \frac{\partial \varepsilon_{\beta\delta}}{\partial r_\gamma}$$

Polarization response to a strain gradient

3 Contributions to $\mu_{\alpha\beta,\gamma\delta}$:

- Electronic (clamped-ion)
- Lattice
- Mixed

Clamped-Ion Flexoelectric tensor

Spatial dispersion of CI piezoelectric tensor

$$e_{\alpha\beta\delta} \propto \left. \frac{d^2 E}{d\varepsilon_\alpha d\eta_{\beta\delta}} \right|_{\mathbf{q}=0} = E^{\varepsilon_\alpha \eta_{\beta\delta}}$$

$$\mu_{\alpha\beta,\gamma\delta} \propto \left. \frac{d^3 E}{d\varepsilon_\alpha d\eta_{\beta\delta} dq_\gamma} \right|_{\mathbf{q}=0} = E_\gamma^{\varepsilon_\alpha \eta_{\beta\delta}}$$

X Electric field and strain perturbations formulated at $\mathbf{q}=0$

$$\varepsilon_\alpha^{\mathbf{q}} \leftarrow \frac{dA_\alpha^{\mathbf{q}}}{dt} \leftarrow \text{Vector potential}$$

$$\eta_{\beta\delta}^{\mathbf{q}} \leftarrow \frac{d(\beta)^{\mathbf{q}}}{dq_\delta} \leftarrow \text{Metric wave}$$

Long-wave DFPT formulation of CI FxE tensor

METRIC WAVE - HOMOGENEOUS STRAIN

$$\mu_{\alpha\beta,\gamma\delta} = \frac{1}{\Omega} E_{\gamma\delta}^{\mathcal{E}_\alpha^*(\beta)}$$

$$\hat{H}_{\mathbf{k},\delta}^{(\beta)} = i\hat{H}_{\mathbf{k}}^{\eta\beta\delta}$$

$$|u_{m\mathbf{k},\delta}^{(\beta)}\rangle = i|u_{m\mathbf{k}}^{\eta\beta\delta}\rangle$$

NEW OBJECTS

$$\tilde{E}_{\gamma\delta}^{\mathcal{E}_\alpha^*(\beta)} = s \int_{\text{BZ}} [d^3k] \sum_m \tilde{E}_{m\mathbf{k},\gamma\delta}^{\mathcal{E}_\alpha^*(\beta)} + \frac{i}{2} \int_{\Omega} \int K_\gamma(\mathbf{r}, \mathbf{r}') n^{\mathcal{E}_\alpha}(\mathbf{r}) n^{\eta\beta\delta}(\mathbf{r}') d^3r d^3r'$$

$$\begin{aligned} \tilde{E}_{m\mathbf{k},\gamma\delta}^{\mathcal{E}_\alpha^*(\beta)} &= i\langle u_{m\mathbf{k}}^{\mathcal{E}_\alpha} | \partial_\gamma \hat{H}_{\mathbf{k}}^{(0)} | u_{m\mathbf{k}}^{\eta\beta\delta} \rangle + i\langle u_{m\mathbf{k}}^{\mathcal{E}_\alpha} | \partial_\gamma \hat{Q}_{\mathbf{k}} \hat{\mathcal{H}}_{\mathbf{k}}^{\eta\beta\delta} | u_{m\mathbf{k}}^{(0)} \rangle + i\langle u_{m\mathbf{k}}^{(0)} | \hat{V}^{\mathcal{E}_\alpha} \partial_\gamma \hat{Q}_{\mathbf{k}} | u_{m\mathbf{k}}^{\eta\beta\delta} \rangle \\ &+ \frac{1}{2} \langle u_{m\mathbf{k}}^{\mathcal{E}_\alpha} | \hat{H}_{\mathbf{k},\gamma\delta}^{(\beta)} | u_{m\mathbf{k}}^{(0)} \rangle + i\langle i u_{m\mathbf{k},\gamma}^{A_\alpha} | u_{m\mathbf{k}}^{\eta\beta\delta} \rangle \end{aligned}$$

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Dynamical quadrupoles

Second moment of the charge response to an atomic displacement

$$Q_{\kappa\beta}^{\mathbf{q}} = \int_{\Omega} \rho_{\mathbf{q}}^{\tau_{\kappa\beta}}(\mathbf{r}) d^3r = -iq_{\beta}Z_{\kappa} + \underbrace{2E_{\mathbf{q}}^{\varphi^*} \tau_{\kappa\beta}}_{d^2E} \frac{1}{d\varphi_{-\mathbf{q}} d\tau_{\kappa\beta,\mathbf{q}}}$$

$$Q_{\kappa\beta}^{\mathbf{q}} = -iq_{\gamma} \underbrace{Q_{\kappa\beta}^{(1,\gamma)}}_{\text{Born effective charge}} - \frac{q_{\gamma}q_{\delta}}{2} \underbrace{Q_{\kappa\beta}^{(2,\gamma\delta)}}_{\text{Quadrupole}} + \dots$$

Born effective charge

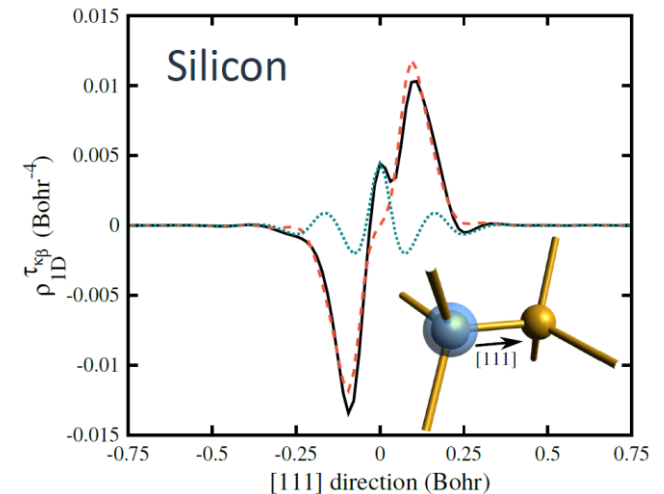
$$\delta_{\beta\gamma}Z_{\kappa} + 2E_{\gamma}^{\varphi^*} \tau_{\kappa\beta}$$

Quadrupole

$$2E_{\gamma\delta}^{\varphi^*} \tau_{\kappa\beta}$$



ONLY AT NON CENTROSYMMETRIC
ATOMIC POSITIONS



Long-wave DFPT formulation of dynamic quadrupoles

$$Q_{\kappa\beta}^{(2,\gamma\delta)} = -2E_{\gamma\delta}^{\varphi* \tau_{\kappa\beta}}$$

SCALAR POTENTIAL - ELECTRIC FIELD

$$|u_{m\mathbf{k}}^{\mathcal{E}_\delta}\rangle = |iu_{m\mathbf{k},\delta}^\varphi\rangle$$

$$E_{\gamma\delta}^{\varphi* \tau_{\kappa\beta}} = -iE_\gamma^{\mathcal{E}_\delta* \tau_{\kappa\beta}} - iE_\delta^{\mathcal{E}_\gamma* \tau_{\kappa\beta}}$$

New Objects

$$E_\gamma^{\mathcal{E}_\delta* \tau_{\kappa\beta}} = s \int_{\text{BZ}} [d^3k] \sum_m E_{m\mathbf{k},\gamma}^{\mathcal{E}_\delta* \tau_{\kappa\beta}} + \frac{1}{2} \int_\Omega \int K_\gamma(\mathbf{r}, \mathbf{r}') n^{\mathcal{E}_\delta}(\mathbf{r}) n^{\tau_{\kappa\beta}}(\mathbf{r}') d^3r d^3r'$$

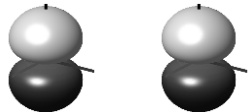
$$E_{m\mathbf{k},\gamma}^{\mathcal{E}_\delta* \tau_{\kappa\beta}} = \langle u_{m\mathbf{k}}^{\mathcal{E}_\delta} | \partial_\gamma \hat{H}_{\mathbf{k}}^{(0)} | u_{m\mathbf{k}}^{\tau_{\kappa\beta}} \rangle + \langle u_{m\mathbf{k}}^{\mathcal{E}_\delta} | \partial_\gamma \hat{Q}_{\mathbf{k}} \hat{\mathcal{H}}_{\mathbf{k}}^{\tau_{\kappa\beta}} | u_{m\mathbf{k}}^{(0)} \rangle + \langle u_{m\mathbf{k}}^{(0)} | V^{\mathcal{E}_\delta} \partial_\gamma \hat{Q}_{\mathbf{k}} | u_{m\mathbf{k}}^{\tau_{\kappa\beta}} \rangle +$$

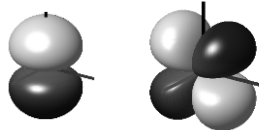
$$\langle u_{m\mathbf{k}}^{\mathcal{E}_\delta} | \hat{H}_{\mathbf{k},\gamma}^{\tau_{\kappa\beta}} | u_{m\mathbf{k}}^{(0)} \rangle + \langle i u_{m\mathbf{k},\gamma}^{A_\delta} | u_{m\mathbf{k}}^{\tau_{\kappa\beta}} \rangle$$

M. Royo and M. Stengel, PRX (accepted)

Why to care about dynamic quadrupoles?

Long-range interatomic forces

$$\Phi_{\kappa\alpha,\kappa'\beta}^{\mathbf{q},DD} = \frac{4\pi}{\Omega} \frac{(\mathbf{q} \cdot \mathbf{Z}_{\kappa}^*)_{\alpha} (\mathbf{q} \cdot \mathbf{Z}_{\kappa'}^*)_{\beta}}{\mathbf{q} \cdot \boldsymbol{\epsilon} \cdot \mathbf{q}},$$

 $\approx d^{-3}$

$$\begin{aligned} \Phi_{\kappa\alpha,\kappa'\beta}^{\mathbf{q},DQ} = & -i \frac{4\pi}{2\Omega} \frac{(\mathbf{q} \cdot \mathbf{Z}_{\kappa}^*)_{\alpha} (\mathbf{q} \mathbf{q} \cdot \mathbf{Q}_{\kappa'}^*)_{\beta}}{\mathbf{q} \cdot \boldsymbol{\epsilon} \cdot \mathbf{q}} \\ & + i \frac{4\pi}{2\Omega} \frac{(\mathbf{q} \mathbf{q} \cdot \mathbf{Q}_{\kappa}^*)_{\alpha} (\mathbf{q} \cdot \mathbf{Z}_{\kappa'}^*)_{\beta}}{\mathbf{q} \cdot \boldsymbol{\epsilon} \cdot \mathbf{q}}, \end{aligned}$$

 $\approx d^{-4}$

Frozen-ion piezoelectric tensor (Martin's theory, 1972)

$$\bar{e}_{\alpha\beta\gamma} = \left. \frac{\partial P_{\alpha}}{\partial \epsilon_{\beta\gamma}} \right|_{\text{FI}} \qquad \bar{e}_{\alpha\beta\gamma} + \bar{e}_{\gamma\beta\alpha} = \frac{1}{\Omega} \sum_{\kappa} Q_{\kappa\beta}^{(2,\alpha\gamma)}$$

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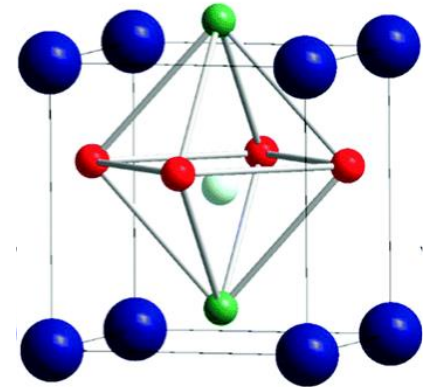
Quadrupoles testcase: Tetragonal PbTiO₃

All calculations are performed using the LDA and norm conserving PSPs

E_{cut}=70 Ha and 8x8x8 MP k-points

	$\kappa = \text{Pb}$	$\kappa = \text{Ti}$	$\kappa = \text{O}_1$	$\kappa = \text{O}_2$	$\kappa = \text{O}_3$
$Q_{\kappa 3}^{(2,11)}$	2.264	-3.545	2.884	-4.186	0.406
$Q_{\kappa 3}^{(2,22)}$	2.264	-3.545	-4.186	2.884	0.406
$Q_{\kappa 1}^{(2,31)}$	-0.062	-3.799	3.123	-1.115	-1.784
$Q_{\kappa 2}^{(2,32)}$	-0.062	-3.799	-1.115	3.123	-1.784
$Q_{\kappa 3}^{(2,33)}$	1.240	-0.195	2.027	2.027	6.653

TABLE I. Quadrupole moments in e-Bohr of PbTiO₃.



Recall: Martin's 1972 formula


$$e_{\alpha\beta\gamma}^P = v_0^{-1} \sum_K \left[\sum_{\delta} e_{K\alpha\delta}^* \Gamma_{K\delta\beta\gamma} - \frac{1}{2} (Q_{K\alpha\beta\gamma} - Q_{K\gamma\alpha\beta} + Q_{K\beta\gamma\alpha}) \right]$$

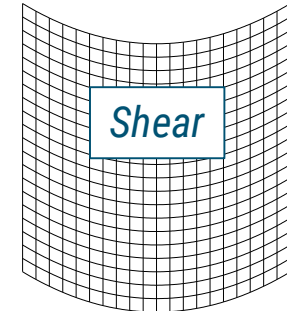
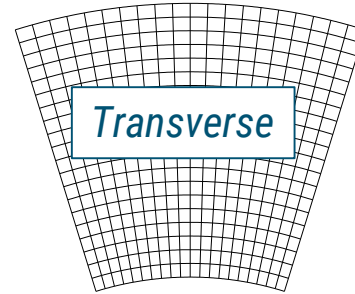
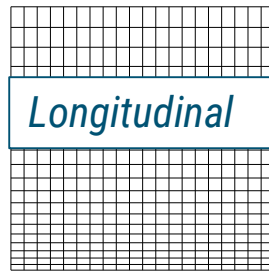
Clamped-ion Piezoelectric Tensor

	$e_{113} = e_{223}$	$e_{311} = e_{322}$	e_{333}
Strain	0.1547	0.3617	-0.8345
Quadrupoles	0.1548	0.3614	-0.8347

TABLE II. Clamped-ion piezoelectric coefficients (in C/m²) of PbTiO₃

Flexoelectric tensor: Cubic materials

Cubic symmetry

 3 independent components



Testcase 1: Isolated noble gas atoms

	μ_L	μ_T	$\mu_S (\times 10^{-4})$
He	-0.479 (-0.479 ^a)	-0.479 (-0.479 ^a)	-0.08 (-0.08 ^a)
Ar	-4.821 (-4.813 ^a)	-4.823 (-4.820 ^a)	-1 (-10 ^a)
Kr	-6.471 (-6.474 ^a)	-6.477 (-6.476 ^a)	-4 (-20 ^a)

TABLE III. Flexoelectric coefficients (pC/m) of noble-gas atom systems. ^a Ref. [3]

() values obtained via numerical derivation in \mathbf{q}
 A. Schiaffino et al. PRB 99, 085107 (2019)

Testcase 2: Real materials

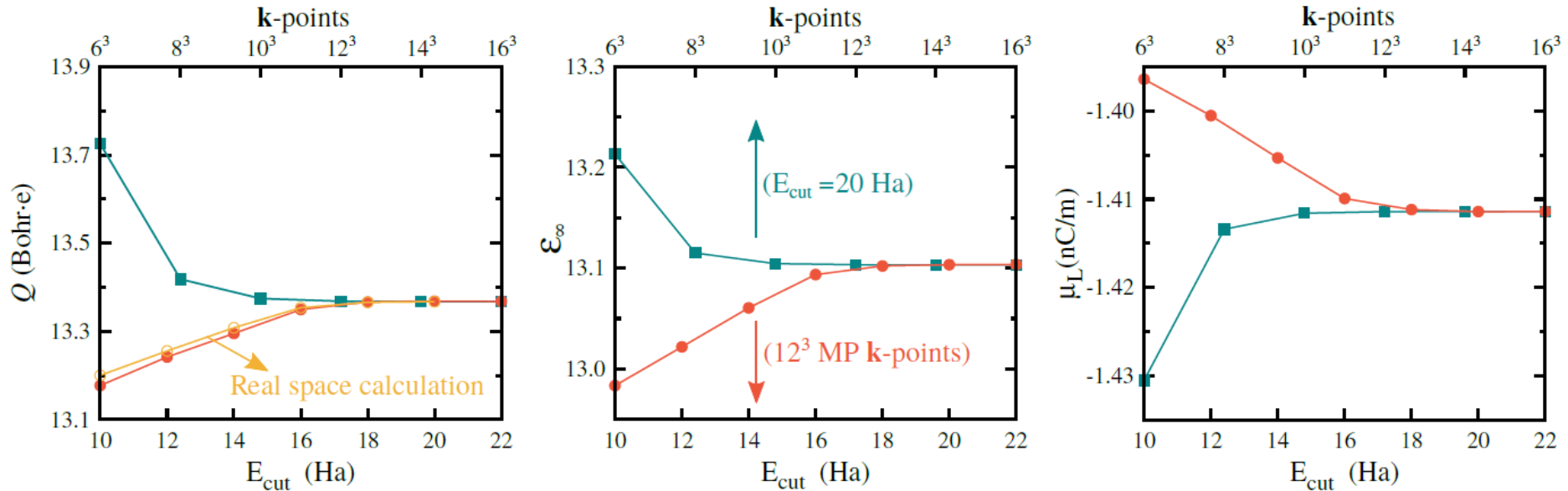
	μ_L	μ_T	μ_S	
A.Schiaffino et al. PRB 99, 085107 (2019)	Si (this work)	-1.4114	-1.0491	-0.1895
	Ref. 3	-1.4110	-1.0493	-0.1894
Stengel PRB 90, 201112(R) (2014)	SrTiO ₃ (this work)	-0.8848	-0.8262	-0.0823
	Ref. 3	-0.8851	-0.8260	-0.0823
	Ref. 6	-0.883	-0.825	-0.082

TABLE V. Flexoelectric coefficients (nC/m) of Si and SrTiO₃.

Convergence study

SYSTEM: Silicon

All calculations are performed using the LDA and norm conserving PSPs



THE SPATIAL-DISPERSION TENSORS CALCULATION REQUIRES A COMPUTATIONAL EFFORT COMPARABLE TO THE CALCULATION OF OTHER STANDARD LINEAR-RESPONSE QUANTITIES

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New objects to implement

$$K_\gamma(\mathbf{G}, \mathbf{G}') = -8\pi G_\gamma \frac{\delta_{\mathbf{G}\mathbf{G}'}}{G^4}$$

hartredq (54_spacepar/m_spacepar.F90)

$$\hat{H}_{\mathbf{k},\gamma}^{\tau_{\kappa\beta}} = V_\gamma^{\text{loc},\tau_{\kappa\beta}} + V_{\mathbf{k},\gamma}^{\text{sep},\tau_{\kappa\beta}}$$

dfpt_vlocaldq (67_common/m_mklocl.F90)

nonlop (choice=22) (66_nonlocal/m_nonlop.F90)

$$\hat{H}_{\mathbf{k},\gamma\delta}^{(\beta)} = \hat{T}_{\mathbf{k},\gamma\delta}^{(\beta)} + V_{\gamma\delta}^{\text{loc},(\beta)} + V_{\mathbf{k},\gamma\delta}^{\text{sep},(\beta)} + \hat{V}_{\delta\gamma}^{\text{H0},(\beta)}$$

mkkin_metdqdq (56_recipspace/m_kg.F90)

dfpt_vlocaldqdq (67_common/m_mklocl.F90)

nonlop (choice=33) (66_nonlocal/m_nonlop.F90)

$$\langle i u_{m\mathbf{k},\gamma}^{A_\alpha} | u_{m\mathbf{k}}^{\eta\beta\delta} \rangle \rightarrow -\frac{i}{2} \underbrace{\langle \tilde{\partial}_{\alpha\gamma} u_{m\mathbf{k}}^{(0)} | u_{m\mathbf{k}}^{\eta\beta\delta} \rangle}_{u_{m\mathbf{k}}^{k_\alpha k_\gamma}} - \frac{i}{2} \underbrace{\langle u_{m\mathbf{k},\alpha\gamma}^{CG} | u_{m\mathbf{k}}^{\eta\beta\delta} \rangle}_{\text{Response to an orbital } \mathbf{B}\text{-field}}$$

NOT IMPLEMENTED

Example of input file

*# Crystalline silicon: computation of the Quadrupole
and CI FxE Tensors*

ndtset 5

#Set 1: Ground state self-consistency

getwfk1 0
kptopt1 1
nqpt1 0
tolvrs1 1.0d-18

#Set 2: Response function calculation of d/dk

iscf2 -3
kptopt2 2
rfelfd2 2
tolwfr2 1.0d-22
rfdir2 1 1 1

#Set 3: Response function calculation of d2/dkdk

getddk3 2
iscf3 -3
kptopt3 2
rf2_dkdk3 1
tolwfr3 1.0d-22

*#Set 4 : Response function calculation of Q=0 phonons,
electric field and strain perturbations*

getddk4 2
kptopt4 2
rfelfd4 3
rfphon4 1
rfatpol4 1 2
rfdir4 1 1 1
tolvrs4 1.0d-10
prepalw4 1 **# Deactivates symmetries for the lw routines**

#Set 5: Long-wave magnitudes calculation

optdriver5 10 # Activates long-wave driver
kptopt5 2
get1wf5 4
get1den5 4
getddk5 2
getdkdk5 3
lw_qdrpl5 1 # Calculate Quadrupoles
lw_flexo5 2 # Calculate CI flexoelectric tensor

#Common input variables

getwfk 1
useylm 1
nqpt 1
qpt 0.0E+00 0.0E+00 0.0E+00
...
...

Example of output files

abi_out

Quadrupole tensor, in cartesian coordinates,

atom	atddir	efidir	qgrdir	real part	imaginary part
1	1	1	1	-0.0000000044	0.0000000000
2	1	1	1	0.0000000044	0.0000000000
1	2	1	1	0.0000000000	0.0000000000
...					
...					
1	2	2	1	-0.0000000021	0.0000000000
2	2	2	1	0.0000000022	0.0000000000
1	3	2	1	13.3682664286	0.0000000000
2	3	2	1	-13.3682664284	0.0000000000
1	1	3	1	-0.0000000023	0.0000000000
...					

Electronic flexoelectric tensor, in cartesian coordinates,

efidir	qgrdir	strdir1	strdir2	real part	imaginary part
1	1	1	1	-0.4661642508	0.0000000000
2	1	1	1	-0.0000000000	0.0000000000
3	1	1	1	-0.0000000000	0.0000000000
1	2	1	1	-0.0000000000	0.0000000000
2	2	1	1	-0.3465045498	0.0000000000
3	2	1	1	0.0000000000	0.0000000000
1	3	1	1	-0.0000000000	0.0000000000
2	3	1	1	0.0000000000	0.0000000000
3	3	1	1	-0.3465045498	0.0000000000
...					

_O_DS5_DDB

**** Database of total energy derivatives ****

Number of data blocks= 1

3rd derivatives - # elements : 216

qpt	0.00000000E+00	0.00000000E+00	0.00000000E+00	1.0	0.00000000E+00	0.00000000E+00	0.00000000E+00	1.0	0.00000000E+00	0.00000000E+00	0.00000000E+00	1.0
1	1	1	1	1	10	0.000000000000000D+00	0.22644265123610D-14					
2	1	1	1	1	10	0.000000000000000D+00	0.24355993235786D+00					
3	1	1	1	1	10	0.000000000000000D+00	0.17956332706089D+00					
1	2	1	1	1	10	0.000000000000000D+00	0.45460672058968D+01					
2	2	1	1	1	10	0.000000000000000D+00	0.24655503777839D+03					
3	2	1	1	1	10	0.000000000000000D+00	0.24726956950978D+03					
1	4	1	1	1	10	0.000000000000000D+00	0.58234827336192D+02					
2	4	1	1	1	10	0.000000000000000D+00	-0.29456056346091D+02					
3	4	1	1	1	10	0.000000000000000D+00	-0.29063615520488D+02					
1	1	2	1	1	10	0.000000000000000D+00	-0.24355993235786D+00					
2	1	2	1	1	10	0.000000000000000D+00	-0.10841168799945D-14					
3	1	2	1	1	10	0.000000000000000D+00	0.26583336681518D+00					
1	2	2	1	1	10	0.000000000000000D+00	0.24656233568811D+03					
2	2	2	1	1	10	0.000000000000000D+00	0.24376369487631D+03					
3	2	2	1	1	10	0.000000000000000D+00	0.24643072491876D+03					
1	4	2	1	1	10	0.000000000000000D+00	0.35339045543433D+00					
....												



ddq ipert=natom+8

State of the implementation

NOT YET MERGED WITH THE TRUNK

Current limitations:

- Perturbations symmetries deactivated
- LDA exclusive
- Not adapted for non-linear core corrections
- $k_{\text{ptopt}} \neq 1$
- $\text{useym} = 1$

TO DO

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Conclusions and Outlook

- CI FxE and quadrupole tensor from a multi-dataset ABINIT run
- No *new* ddq response functions required
- Little computational cost
- Developing of full FxE tensor (lattice and mixed contribs.)
- Other spatial dispersion properties (natural optical/acoustical activity)

THANK YOU!

M. Royo and M. Stengel, Phys. Rev. X (accepted), arXiv:1812.05935