

The Multibinit project





AbiDev 2019

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### Motivations

# Understanding and engineering functional properties often require undestanding materials at the atomic scale.



- From *ab-initio* data to second-principles.
- Integrate some degress of freedom.
- Access relevant properties at operating conditions from different length-scales.



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### Multibinit in the package

Multibinit

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- Multibinit is a new executable included in the *src*/98\_*main/multibinit*
- New directory *src*/78\_*effpot* have been created with many new files !
  - New directory structure is planned to better organize the code
  - Split abastract layer from each potential
  - Split the mover from the potentials
- New input have been created (mix between anaddb and ABINIT)
- The parsing of the XML file is done with Fortran or LibXML (more efficient...)
- MPI Parallelization
- Hybrid Monte Carlo is implemented
- Some of abinit movers are used
- New automatic tests in the version 8
- New documentation with a topic on the website.
- New tutorial to learn how to use multibinit (lattice and spin)

2



#### Basic Procedure

- Express energy as low order Taylor expansion
  - Use hight-symmetry reference structure
  - Variables: local modes and strains
- Determine coefficients from a limited number of DFT calculations
- Resulting energy expansion allows for finite temperature Monte Carlo

Successful results already for many oxide systems. However identifying a small set of degrees of fredom is often difficult.

# We need a general description of the energetics of all the atomic degress of freedom

K. Rabe and J. Joannopoulos, PRL 59, 570 (1987); PRB 36, 6631 (1987)
 W. Zhong et al., PRL 73, 1861 (1994); PRB 52, 6301 (1995)



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Energy changes around reference structure due to distortions

 $E_{eff}(\{\mathbf{u}_i\},\eta) = E_{\rho}(\{\mathbf{u}_i\}) + E_{s}(\eta) + E_{s\rho}(\{\mathbf{u}_i\},\eta)$ 

J. Wojdel et al., J. Phys. Condens. Matter 25 (2013) 305401

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Energy changes around reference structure due to distortions

$$E_{eff}(\{\mathbf{u}_i\}, \eta) = E_p(\{\mathbf{u}_i\}) + E_s(\eta) + E_{sp}(\{\mathbf{u}_i\}, \eta)$$
Energy change due to  
strain only
$$E_s^{Harm}(\eta)$$

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J. Wojdel et al., J. Phys. Condens. Matter 25 (2013) 305401



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Energy changes around reference structure due to distortions



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J. Wojdel et al., J. Phys. Condens. Matter 25 (2013) 305401



Energy changes around reference structure due to distortions



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Energy changes around reference structure due to distortions



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Energy changes around reference structure due to distortions



Same framework as  $H_{eff}$  but includes all-atomic degrees of freedom. J. Wojdel et al., J. Phys. Condens. Matter 25 (2013) 305401

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**CESAM** 3. Lattice Effective Hamiltonians



Phonon term 
$$E_p({\mathbf{u}_i}) = E_p^{Harm}({\mathbf{u}_i}) + E_p^{Anharm}({\mathbf{u}_i})$$

Taylor development around the reference structure

$$E_{p}(\{\mathbf{u}_{i}\}) = \underbrace{\frac{1}{2} \sum_{i \alpha j \beta} \mathcal{K}_{i \alpha j \beta}^{(2)} u_{i \alpha} u_{j \beta}}_{\text{Harmonic terms}} + \underbrace{\frac{1}{6} \sum_{i \alpha j \beta k \gamma} \mathcal{K}_{i \alpha j \beta k \gamma}^{(3)} u_{i \alpha} u_{j \beta} u_{k \gamma} + \mathcal{O}(u^{4})}_{\text{Anharmonic terms}}$$

with  $K_{i\alpha j\beta \dots}^n = \left. \frac{\partial^n E_{eff}}{\partial u_{i\alpha} \partial u_{j\beta} \dots} \right|_{\eta=0}$ .

- Must comply the Acoustic Sum Rule ightarrow difficulte to enforce for n>2
- Use displacement differences  $\rightarrow$  ASR satisfied by construction

$$E_{\mathsf{Harm}}\left(\{u_i\}
ight) = rac{1}{2}\sum_{\substack{i,j,k,h\lpha,eta}} ilde{\kappa}_{ijlpha kheta}^{(2)}(u_{ilpha} - u_{jlpha})(u_{keta} - u_{heta})$$

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$$E_{p}^{Harm} \left( \left\{ \mathbf{u}_{i} \right\} \right) = E_{p,SR}^{Harm} \left( \left\{ \mathbf{u}_{i} \right\} \right) + E_{p,LR}^{Harm} \left\{ \left\{ \mathbf{u}_{i} \right\} \right)$$

$$E_{p}^{Harm} \left( \left\{ \mathbf{u}_{i} \right\} \right) = \frac{1}{2} \sum_{i\alpha j\beta} K_{i\alpha j\beta}^{(2)} u_{i\alpha} u_{j\beta} = \frac{1}{2} \sum_{i\alpha j\beta} \left( S_{i\alpha j\beta} + L_{i\alpha j\beta} \right) u_{i\alpha} u_{j\beta}$$

$$= \frac{1}{2} \sum_{i \in SC} u_{i\alpha} \sum_{j \in SR} S_{i\alpha j\beta} u_{j\beta} + \frac{1}{2} \sum_{i \in SC} u_{i\alpha} \sum_{j \in SC} \left( \sum_{\beta} L_{i\alpha j\beta}(b) \right) \underbrace{u_{j\beta}^{b}}_{L_{i\alpha j\beta}(\mathbf{q}=\mathbf{0})}$$

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### $E_p^{Anharm}\left(\left\{\mathbf{u}_i\right\}\right)$

- Anharmonicity degree : chose from maximum order: 3, 4, 5, 6,...
- For each order build the possible terms

$$E_{\rho}^{(n)}(\{\mathbf{u}_{i}\}) = \frac{1}{!n} \sum_{\substack{i,j,k,l,m,n,\cdots \\ \alpha,\beta,\gamma,\cdots}} \mathcal{K}_{ijklmn}^{(n)\alpha\beta\gamma}(u_{i\alpha} - u_{j\alpha})(u_{k\beta} - u_{l\beta})(u_{m\gamma} - u_{n\gamma})\cdots$$

- ASR enforces by using differences
- Too many terms for "hight" orders
- Use symmetries to reduce de number of coefficients



 $E_p^{Anharm}(\{\mathbf{u}_i\})$ Symmetry considerations: "Symmetry adapted terms"

The operations of the space group can generate a collection of symmetry related terms which reduces the number of *independent* parameters. In practice :

- Define range or cutoff radius
- Build all terms within this range
- Apply all symmetries to find relations between terms :



 $(u_{0Bx} - u_{0O3x})^2 (u_{0By} - u_{0O3y})^2 + \frac{\pi}{2}$  rotation about the y-axis transforms to  $(u_{0Bz} - u_{0O1z})^2 (u_{0By} - u_{0O1y})^2 \rightarrow 15$  independent parameters (1<sup>st</sup> nearest neighbors at 3<sup>rd</sup> and 4<sup>th</sup> order in *u*) for the anharmonic phonon part.

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## Phonon term $E_p({\mathbf{u}_i}) = E_p^{Harm}({\mathbf{u}_i}) + E_p^{Anharm}({\mathbf{u}_i})$

#### Harmonic terms

- $K^{(2)}$  obtained by DFPT as the back-Fourier-transformed dynamical matrix
- Remove dipole-dipole interaction in real space to obtain the SR only.
- Use ewald summation to compute the exact dipole-dipole interaction.
- $\Rightarrow$  Harmonic terms exact by construction and include all phonon branches



## Phonon term $E_p({\mathbf{u}_i}) = E_p^{Harm}({\mathbf{u}_i}) + E_p^{Anharm}({\mathbf{u}_i})$

#### Harmonic terms

- $K^{(2)}$  obtained by DFPT as the back-Fourier-transformed dynamical matrix
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#### Anharmonic terms

- Displacement differences expression allows for an unconstrained optimization of  $\tilde{K}^{(n)}$
- Determined by fitting key quantites to limited number of DFT calculations (training set)



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### Strain term $E_{s}(\eta)$

Taylor development around the reference structure

$$E_{s}(\eta) = \frac{N}{2} \sum_{ab} C_{ab}^{(2)} \eta_{a} \eta_{b} + \frac{N}{6} \sum_{abc} C_{abx}^{(3)} \eta_{a} \eta_{b} \eta_{c} + \mathcal{O}(\eta^{4})$$
with  $C_{ab...}^{(m)} = \frac{1}{N} \frac{\partial^{m} E_{eff}}{\partial \eta_{a} \partial \eta_{b} \cdots} \Big|_{u_{i}=0}$ 

• Harmonic term calculated directly : the frozen ion elastic constants (DFPT)

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• Anharmonic terms : in practice not required for semi-quantitative results

#### J. Wojdel et al., J. Phys. Condens. Matter 25 (2013) 305401



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#### LIÈGE université 3. Lattice Effective Hamiltonians

## Strain-Phonon coupling $E_{sp}(\{u_i\},\eta)$

Taylor development also subject to ASR

$$E_{sp}\left(\left\{ u_{i}
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ight\right)$$





Forces at homogeneous strain

Change in force csts with strain

- The  $\Lambda^{(m,n)}$  must comply with a set of ASRs
- Use displacement differences

$$\begin{split} E_{sp}\left(\left\{u_{i}\right\},\eta\right) &= \frac{1}{2}\sum_{a}\sum_{ij\alpha}\Lambda_{ai\alpha}^{(1,1)}\eta_{a}\left(u_{i\alpha}-u_{j\alpha}\right) \\ &+ \frac{1}{6}\sum_{a}\sum_{ij\alpha kh\beta}\Lambda_{ai\alpha j\beta}^{(1,2)}\eta_{a}\left(u_{i\alpha}-u_{j\alpha}\right)\left(u_{k\beta}-u_{h\beta}\right)+\dots \end{split}$$

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#### Fitting anhamonic terms

Fit is performed on both Phonon and Strain-phonon terms.

• For a set (*TS*) of configurations (*s*) from DFT calculations, we fit with least squares method, the goal function is defined as:

$$G[\lambda_{\rho}, TS] = \frac{1}{M_1} \sum_{s,\alpha,j} \left( F_{\alpha_j}^{TS}(s) - F_{\alpha_j}[\lambda_{\rho}](s) \right)^2 + \frac{1}{M_2} \sum_{s,j} \Omega^2(s) \left( \sigma_j^{TS}(s) - \sigma_j[\lambda_{\rho}](s) \right)^2$$

• where 
$$\Omega(s) = ig(V(s)\sqrt(N)ig)^{(-1/3)}$$

- The goal function has to satisfy  $\frac{\partial G[\lambda_p, TS]}{\partial \lambda_\mu} = 0 \quad \forall \mu \text{ and } \frac{\partial^2 G[\lambda_p, TS]}{\partial \lambda_\mu \partial \lambda_\nu} \ge 0 \quad \forall \mu \nu$
- We solve the system of p linear equations in order to get the set of coefficients  $\lambda_p$

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```
prt_model = 1
#
#Inputs for the fitted coefficients
# _ _ _ _
fit_coeff = 1
fit_generateCoeff = 1
fit_rangePower = 3 4
                        # Range for the powers of the polynomial (default 3
   to 4 )
fit_ncoeff = 7
                        # Number of coefficients to consider for the fit
fit_cutoff = 8
                        # Cutoff for the interactions
fit_anhaStrain = 0
fit_SPCoupling = 1
ts_option = 1
```



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#### Example of CaTiO<sub>3</sub> cubic phase



• Mean Standard Deviation of the Energy with only Harmonic: 250.58 meV/f.u.

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#### Example of CaTiO<sub>3</sub> cubic phase



• Mean Standard Deviation of the Energy with only Harmonic: 250.58 meV/f.u.

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• Mean Standard Deviation of the Energy with 1 Coeff: 284.70 meV/f.u.



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#### Example of CaTiO<sub>3</sub> cubic phase



• Mean Standard Deviation of the Energy with only Harmonic: 250.58 meV/f.u.

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• Mean Standard Deviation of the Energy with 5 Coeffs: 21.28 meV/f.u.



#### Example of CaTiO<sub>3</sub> cubic phase



• Mean Standard Deviation of the Energy with only Harmonic: 250.58 meV/f.u.

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• Mean Standard Deviation of the Energy with 10 Coeffs: 12.75 meV/f.u.



#### Example of CaTiO<sub>3</sub> cubic phase



• Mean Standard Deviation of the Energy with only Harmonic: 250.58 meV/f.u.

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• Mean Standard Deviation of the Energy with 20 Coeffs: 19.37 meV/f.u.



#### Example of CaTiO<sub>3</sub> cubic phase



• Mean Standard Deviation of the Energy with only Harmonic: 250.58 meV/f.u.

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• Mean Standard Deviation of the Energy with 50 Coeffs: 19.03 meV/f.u.



#### Example of CaTiO<sub>3</sub> cubic phase



• Mean Standard Deviation of the Energy with only Harmonic: 250.58 meV/f.u.

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• Mean Standard Deviation of the Energy with 100 Coeffs: 6.46 meV/f.u.



#### Example of CaTiO<sub>3</sub> cubic phase



• Mean Standard Deviation of the Energy with only Harmonic: 250.58 meV/f.u.

• Mean Standard Deviation of the Energy with 175 Coeffs: 4.45 meV/f.u.



#### Example of CaTiO<sub>3</sub> cubic phase



• Mean Standard Deviation of the Energy with only Harmonic: 250.58 meV/f.u.

• Mean Standard Deviation of the Energy with 175 Coeffs: 4.45 meV/f.u.



#### Example of CaTiO<sub>3</sub> cubic phase



• Mean Standard Deviation of the Energy with only Harmonic: 250.58 meV/f.u.

• Mean Standard Deviation of the Energy with 175 Coeffs: 4.45 meV/f.u.



Bounding the potential

- If the highest order is odd
- If the highest order is even and the coefficient negative
- $\Rightarrow$  The potential might be unstable !!
  - Add artifical physics with maths to bound the process.
- $\rightarrow$  See M. Schmitt presentation.





# Analyzing the results $_{q(Agate)}$

- Visualize the trajectory
- Perform all "MD" analysis (PDF, T, P, V, ...)
- Extract phonons at finite temperature (with a-TDEP see F. Bottin talk)
- Project on modes and follow phase transitions

	qTdep			- 0 ×			
t37_HIST.nc			qTdep		_ 0 ×		
Supercell	t37_HIST.nc				qTdep	- 0 )	
Trajectory (101 step(s))	Supercell		t37_HIST.nc		Band Structure		
First time: 0	Unit cell		Supercell				
Last time 100 🗘	Lattice		Unit cell		600	-	
Step 1 0	Lattice scaling 14.852957 14.852957 14.852957		Options				
Temperature 495,05K	a 0.5 0	0	Order expansion 2 *		500		
Multiplicity	b 0 0.5	0	radius cutoff for order 2 7.426	48 bohr +	400		
2 0 0	c 0 0	0.5			1-		
0 2 0	Space group: 221: Pm-3m		DOS smearing 4.5er	6 Ha +	300		
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			kpt=0.177122 E=717.256				

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#### Lattice part

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- Multibinit is able to make interface between first principles calculations and mesoscopic simulations for «any» system
- Automatic construction of the models with harmonic and anharmonic contributions from first principles
- Automatic bound process

6. Conclusion

- Tool able to run dynamics (Monte Carlo or Molecular dynamics)
- Good tools for result analysis and post-processing of data ((q)Agate)
- New ionmov in abinit to generate the training set



#### Miscellanous

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- Spin potential available (X. He talk)
- Spin dynamics available (X. He Talk)
- Spin-lattice coupling (N. Helbig)

Conclusion

- Effective Hamiltonian WIP (W. Lafargue-dit-Hauret poster)
- Coupling with electronic potential (M. Schmitt and previous talks)
- Common data structure for ease of use and interfacing







### Thank You !

- ULiège : Marcus Schmitt, He Xu, Nicole Helbig, Fabio Ricci, Matthieu Verstraete, Eric Bousquet, Philippe Ghosez
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Interactions also with :

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- LIST : Jorge Íñiguez

